

Structural, Dielectric and Ferroelectric Properties of Tungsten Substituted Barium Titanate Ceramics

SHEELA DEVI and A. K. JHA*

*Thin Film & Materials Science Laboratory, Delhi College of Engineering, (Faculty of Technology, University of Delhi), New Delhi - 42,
E-mail: dr_jha_ak@yahoo.co.in.*

In the present work, X-ray diffractograms analysis reveals the formation of the single phase crystalline structure. Tetragonal strain of the samples have been calculated and it shows a decrease with increasing sintering temperature. The variation of dielectric constant (ϵ_r) with temperature shows that the samples undergo a diffuse type ferro - paraelectric phase transition. The dielectric constant is found to increase with increasing sintering temperature. Variation of dielectric loss ($\tan\delta$) with temperature have also been measured. Further, the diffusivity of the samples has been calculated and is reported in this paper. The P-E studies of the W-substituted barium titanate ceramics with temperatures are reported.

Key Words: Dielectric Response, Ferroelectrics, Solid-State Reaction, X-Ray Diffraction.

INTRODUCTION

Since the discovery of barium titanate (BaTiO_3) in early 1940s, this material has been used for a wide range of scientific and industrial applications such as capacitors, ultrasonic transducers, pyroelectric infrared sensors, positive temperature coefficient (PTC) resistors¹⁻⁴, etc. Being a lead free ferroelectric material, BaTiO_3 is an environmental friendly material, making it a good substitute for Pb containing compounds for various applications. It is a typical ferroelectric material with Curie temperature in the range of 120-130 °C. It has a perovskite (ABO_3) tetragonal structure at room temperature⁵. Beyond T_c , in paraelectric phase, its structure transforms from tetragonal to cubic. In cubic structure, it shows positive temperature coefficient of resistance (PTCR) behaviour and is used as material for PTCR devices. The perovskite structure is one of the most versatile structures for tailoring the properties of materials⁶⁻⁸. Both A and B-sites dopants are used to modify the electrical and dielectric properties of BaTiO_3 . A lot of works have been reported on A-site doping⁹⁻¹¹. It is well established that dielectric properties of BaTiO_3 ceramics depend not only on the nature and amount of dopants, but also on the different processing factors¹². Not much work, however, have been reported on B-site doping and there is hardly any extensive report on the tungsten doping in barium titanate. Tungsten has been found to be effective in enhancing dielectric properties of

other ferroelectric materials¹³. In the present work, structural, dielectric and ferroelectric properties of tungsten substituted barium titanate ferroelectric ceramics have been investigated and is reported here.

EXPERIMENTAL

Samples of compositions $\text{Ba}(\text{Ti}_{1-x}\text{W}_x)\text{O}_3$; $x = 0.0$ and 0.15 , were prepared using solid-state reaction method taking BaCO_3 , TiO_2 and WO_3 (all from Aldrich of 99.9% purity) in stoichiometric proportions. The powders were thoroughly mixed, ground and passed through sieve of appropriate size. Mixtures were calcined at 1150°C for 2h in air. The calcined specimens were mixed with appropriate quantity of polyvinyl alcohol and molded into disc shape pellets by applying a pressure of 300MPa. Pellets were sintered at 1200°C and 1300°C for 2h in air. X-ray diffractograms of all the calcined and sintered samples were recorded on a Philips X-ray diffractometer using CuK_α radiations in the range $10^\circ \leq 2\theta \leq 70^\circ$ at a scanning rate of $0.02^\circ/\text{s}$. The sintered pellets were polished to a thickness of 1mm, coated with silver paste on both sides for the use as electrodes and finally cured at 300°C for 15 min. The dielectric parameters were measured from 100 Hz to 1 MHz on an Agilent 4284 LCR meter at oscillation amplitude of 1 volt. P-E hysteresis measurements were done from room temperature to Curie temperature using an automatic P-E loop tracer based on Sawyer – Tower circuit.

RESULTS AND DISCUSSIONS

Fig.1 shows the observed X-ray diffractograms of the studied samples. It is observed that single phase perovskite structure is formed in all the samples. Lattice parameters of the sample were calculated from the observed d-values in the X- ray diffractograms and refined using least square refinement method by a computer programe package – PowderX¹⁴. In the doped sample, some additional peaks are formed which are possibly due to unreacted tungsten oxide.

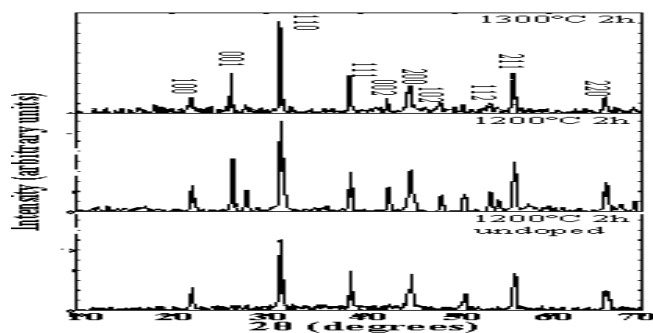


Fig. 1 XRD pattern of the studied sample

On the basis of ionic radii and co-ordination number¹⁵, tungsten is expected to substitute titanium (B) sites (Table 1). It is observed that the lattice parameters and unit cell volume decrease on tungsten substitution (Table 1). The decreased cell volume and lattice parameters implies that the substituted W ion exists in a valence of +6, as the size of W^{6+} is smaller than Ti^{4+} .

A small variation in the position and intensity of the peaks is also observed which can be understood in terms the change in lattice parameters. On substitution of tungsten it is found that the sample is mainly composed of tetragonal phase due to splitting of (002) and (200), (102) and (201), and (112) and (211) occurring at $2\theta = 43.97^\circ$ and 44.83° , 49.69° and 50.33° and 54.95° and 55.45° , respectively¹⁶.

TABLE-1
LATTICE PARAMETERS, IONIC RADII, VOLUME, DIFFUSIVITY AND POLARIZATION VALUE

Temp.	a(A°)	c(A°)	c/a	Ionic radii (A°)	Volume (A°) ³	Diffusivity	2P _r (μC/cm ²)
1200°C 2h Tungsten free	3.999	4.038	1.009	Ti ⁴⁺ = 0.68	64.57	1.18	4.00
1200°C 2h	3.996	4.031	1.008	W ⁶⁺ = 0.60	64.36	1.65	3.24
1300°C 2h	3.986	4.012	1.005	W ⁶⁺ = 0.60	63.74	1.74	8.41

Fig.2 shows the variation of dielectric constant with temperature of the studied samples, measured at different frequencies. It is observed that on tungsten substitution, T_c decreases from 115°C (for $x = 0.0$) to 85°C (for $x = 0.15$). The decrease in Curie temperature can be understood as follows. With tungsten substitution there is a reduction in Curie temperature because the bond strength between W-O is weaker as compared to Ti-O bond¹⁷. This decrease in bond energy leads to weaker distortion of the octahedron, which also decreases c/a ratio (also observed in Table1), resulting in the decrease of Curie temperature. It is observed that the dielectric constant of the samples show a decrease as frequency is increased from 100 Hz to 100 kHz. The dielectric constant (ϵ) of a material has four polarization contributions: electronic polarization (ϵ_e); ionic polarization (ϵ_i); dipolar polarization (ϵ_d) and space charge polarization (ϵ_s)¹⁸. Response frequencies for electronic and ionic polarization are $\sim 10^{16}$ and 10^{13} Hz respectively; and at frequencies above 100 kHz, contribution from space charge polarization is not expected¹⁸. This saturation of space-charge polarization results in low dielectric constant values at higher frequencies. As frequency is increased, dielectric constant is observed to decrease in all the samples, which is a normal

behavior of these materials¹⁹ and at a frequency of 100 kHz, contribution from space charge polarization is likely to be dominant²⁰. The decrease in the dielectric constant is rapid at lower frequencies and become slower at higher frequencies. At the higher frequencies the dipoles can not follow the applied ac electric field resulting in the decrease of dielectric constant at higher frequencies. The high values of dielectric constant at low frequencies can thus be attributed to the presence of such space charges in the samples.

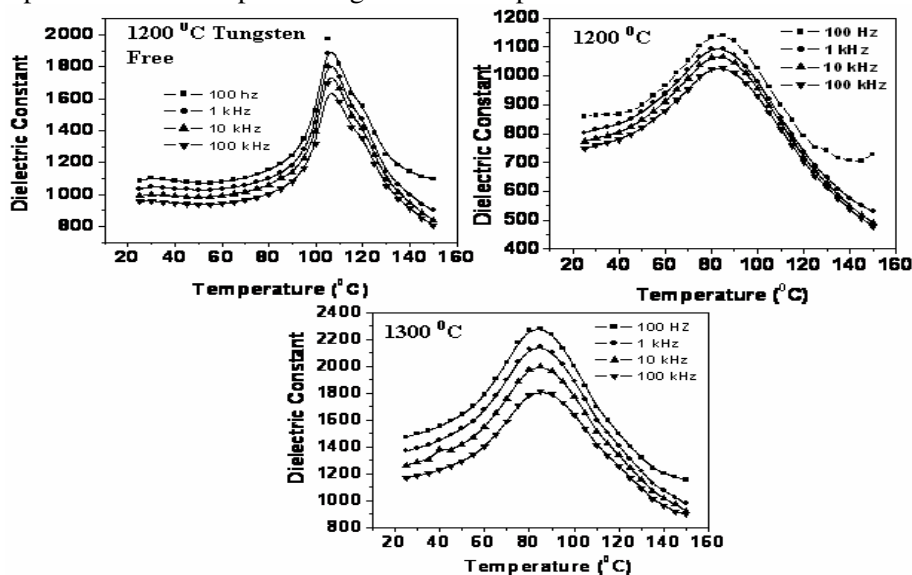
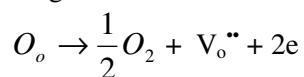


Fig. 2 Dielectric constant vs. temperature at different frequency (a) undoped sample (b) 1200 °C 2h (c) 1300 °C 2h

It is also observed that on increasing sintering temperature the dielectric constant increases. The observed maximum dielectric constant (ϵ_{\max}) of all the studied samples is given in Table 1. The increase in dielectric constant on sintering temperature can be understood as follows. It is known that in barium titanate oxygen vacancies are major structural defects²¹⁻²³. These vacancies are generated due to the loss of oxygen during sintering at high temperature in accordance with the following relation²⁴



The above equation is in accordance with Kröger – Vink notation of defects²⁵. These vacancies act as singly ionized ion at room temperature²⁶. These defects act as space charge (in the form of oxygen vacancies) in the specimen²⁷. In the sample, sintered at higher temperature, larger numbers of oxygen vacancies are created. The increase in dielectric constant value is known to be related to the oxygen ions or oxygen vacancies created during sintering²⁸. This is the possible cause for the increase in dielectric constant in the material.

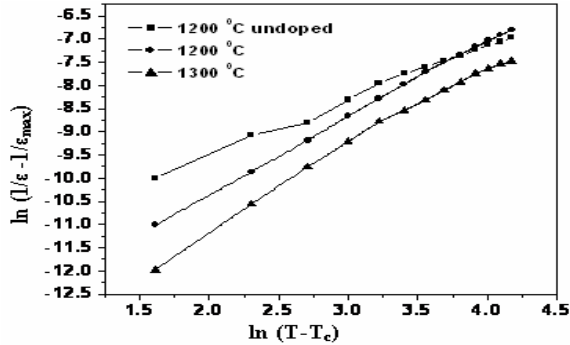


Fig. 3 Variation of $\ln (1/ \varepsilon - 1/ \varepsilon_{\max})$ as a function of $\ln (T- T_c)$.

Fig. 3 shows the variation of $\ln (1/ \varepsilon - 1/ \varepsilon_{\max})$ as a function of $\ln (T- T_c)$. Diffusivity in the sample can be calculated using the expression²⁹

$$\ln (1/ \varepsilon - 1/ \varepsilon_{\max}) = \gamma \ln (T- T_c) + \text{constant} \tag{1}$$

Where ε_{\max} is the maximum dielectric constant, γ is the diffusivity or disorderness of ferroelectric to paraelectric phase transition and its value lies between 1 and 2. When $\gamma = 1$, the material follows the ideal Curie- Weiss law as in the case of normal ferroelectric while when $\gamma = 2$ the material corresponds to a complete diffuse phase transition (DPT)³⁰. The calculated values of diffusivity are given in Table 1. The value of diffusivity increase with increasing sintering temperature indicating increased disorderness in the system. This could be due to the increased defects (here oxygen vacancy) induced disorderness³¹.

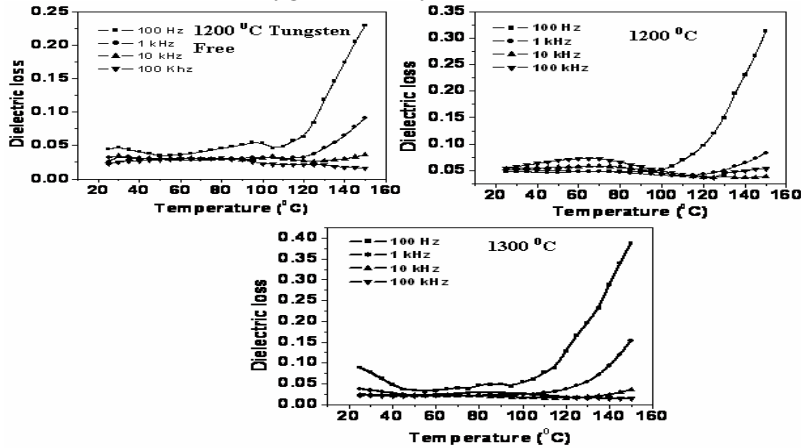


Fig. 4 Dielectric loss vs. temperature at different frequency (a) undoped sample (b) 1200 °C 2h (c) 1300 °C 2h

Fig. 4 shows the variation of dielectric loss with temperature at different frequencies. It shows negligible change upto around 100 °C, largely independent of temperature, thereafter, it increases sharply. The loss is observed to be higher at lower frequency. This can be well known that as the frequency is increased

more and more dipoles are unable to follow the oscillating field resulting in the reduction in loss with increasing frequency. Further, the dielectric loss observed to increase on tungsten substitution, which is possibly due to the generation of space charges in the form oxygen vacancies, as discussed earlier. These space charges are dominant at lower frequencies and higher temperatures³².

Fig. 5 shows the P-E loops of the studied samples at different temperatures. Well-defined hysteresis loops are formed. It is observed that $2P_r$ value increases with increasing sintering temperature. It is also observed on with increasing temperature from room temperature to Curie temperature; polarization decreases obtained at Curie temperature which indicating that the samples become paraelectric at Curie temperature and no loop is formed.

It is known that ferroelectric properties are affected by the composition, microstructure and lattice defects like oxygen vacancies in the structure of materials^{33, 34-35}. In soft ferroelectrics, with higher-valent substituents, the defects are cation vacancies. Higher valent substituents and associated cation vacancies tend to form dipolar defects and random field around these dipolar defects lowers the activation barrier required for nucleation of new domains, leading to lower E_c ^{36, 37}. Also it has been reported that cation vacancies generated by donor doping make domain motion easier and enhance the ferroelectric properties³⁸.

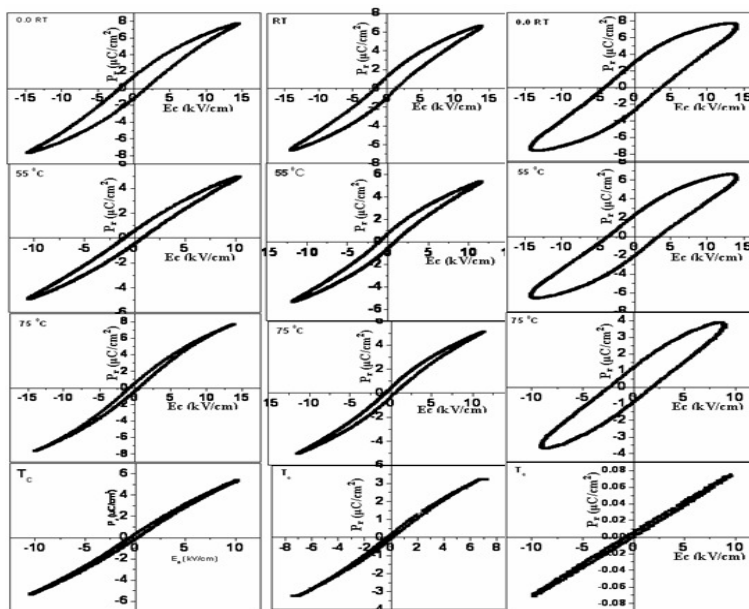


Fig. 5 P-E Hysteresis loop of (a) undoped sample (b) 1200 °C 2h (c) 1300 °C 2h

In the doped samples, the associated cation vacancies formation due to the substitution of Ti^{4+} by W^{6+} suppresses the concentration of oxygen vacancies. A reduction in the number of oxygen vacancies reduces the pinning effect on the domain walls, leading to enhanced remnant polarization³⁹.

Conclusion

X-ray diffractograms of the samples reveal the formation of single phase layered perovskite structure. The transition from ferro to para-electric becomes broader on increasing the sintering temperature as well as on substitution of tungsten indicating diffusivity in the system. Dielectric constant and loss decreases on increasing frequency. Tungsten substitution leads to higher remnant polarization ($2P_r$) values. The P-E loops become thinner as the measurement temperature approaches the Curie temperature.

REFERENCES

1. G.H. Haertling, *J. Amer. Ceram.Soc.* **82**, 797 (1999).
2. R.E.Jones, P.D.Mainar, J.O.Olowolafe, J.O.Campbell AND C. J. Mogab, *Appl. Phys. Lett.* **60**, 1022 (1992).
3. M.H.Frey and D.A.Payne, *ibid.*, **63**, 2753 (1963).
4. H. Hu. And S.B. Krupanidhi, *J. Appl. Phys.* **74**, 3373 (1993).
5. H. D Megaw, *Nature*, **154**; 484, (1954).
6. Z. He, J. Ma, Y.Qu, X .Feng, *J. Eur. Ceram. Soc.*, **22**, 2143 (2002).
7. J. Zhao, Z. Li, Z. Guisens. *Actuators A Phys.* **95**, 46 (2001).
8. J.G. Kim, W-S Cho, K, Park, *Mater.Sci.EngB.* **83**, 123 (2001)
9. R. Farhi, M.El Marssi, A, Simon, and J. Ravez, *J. Eur. Phys.* **B18**, 605- 610 (2000).
10. Zhi- Cheng Li, Hong Zhang, Xiaodong Zou, Bill Bergman, *Materials science & Engineering B*, **116**, 34-39 (2005).
11. Jianquan Qi, Longtu Li, Yongli Wang, Yuwei Fan, Zhilun Gui, *Materials Chemistry and Physics*, **82**, 423-427 (2003).
12. Haifeng, Yinyin Lin, Ting-Ao-Tang, *Solid State Communication.*, **135**, 304 , (2005).
13. I.Coondoo, A. K. Jha, S.K.Aggarwal, and N.C.Soni: *Journal of European Ceramics Society*, **27**, 253-260 (2007)
14. Dong C, PowderX: Window-95 based program for powder X-ray Diffraction data Processing. *J. Appl Cryst.* **32**, 838 (1999).
15. R D Shannon and C T Prewitt, *Acta Cryst B*; **25**, 925 (1969).s
16. Soon- gyu Kwon, Kyoon Choi, Byung-Ik Kim, *Materials Letters*; **60**,979-982 (2006).
17. Cao Wanqiang, Junwen Xiong, Sun Juanpin: *Materials Chemistry and Physics* **106**, 338-342 (2007).
18. Relva C. Buchanan, *Ceramic Materials for Electronic Processing, Properties and Applications*, Marcel Dekker Inc., New York, (1986).
19. Lines M E, *Glass AM, Principle & Applications of ferroelectric materials*, Clarendon Press, Oxford, 102-104
20. H. R Rukmini, R.N.P Choudhary, D.L Prabhakara, *Materials Chemistry and Physics* **64**, 171- 178 (2000).
21. Om Prakash, Devendra Kumar, R. K Dwivedi. K Srivastava, Prakash Singh, Sindhu Singh, J, *Mater Sci*, **42**, 5490-5496 (2007).
22. J.Q. Qi, W.P. Chen, Y.J. Wu, L.T. Li, *J.Am. Ceram. Soc.* **81**, 2, 437-438 (1998).
23. Bin Cui, Pengfei Yu, Jing Tian, Huilin Guo, Zhuguo chang, *Material Science and Engineering A*, **454-455**, 667-672 (2007).
24. Burn I, Neirman, *J. Material Science*, **17**, 3510-3524 (1982).
25. F. A Kroger, H. J Wink: *Solid state Physics*, Academic press, New York, (1956).
26. N. Hirose, J. M. S Skakle, A.R West, *Journal of Electroceramics*, **3**, 233-238 (1999).
27. R. Moos, K. H Hardlt, *J. Amer. Ceramic Society*, **80**, 10, 2549-2562 (1997).
28. Prasun Ganguly, A.K Jha, K.L Deori., *J. Electroceramics* (In press), DOI 10.1007/s10832-007-9337-6

29. S. M Pilgrim, A. E Sutherland, S.R Winzer., *J. Am. Cer. Soc.*, **73**, 10, 3122- 3125 (1990).
30. J. Daniels, K. H Harlet, R. Wernicke, *Philips Res. Rep.* **36**, 487-559 (1976).
31. V. Raghavan, Materials Science and Engineering, *Prentice- Hall of India*, New Delhi 194 (2004).
32. R.C. Buchanan, Ceramic materials for Electronics: Processing, properties and Applications (Marcel Dekker Inc., New York) 38, (1986).
33. H. Watanabe, T. Mihara, H. Yoshimori and C. A. paz. De Arauja, *Jpn. J. Appl. Phys.*, **34**, 5240 (1995).
34. Wei Wang, Jun Zhu, Xiang Yu Mao, Xiao-Bing Chen, *Materials Research Bulletin* **42**, 274-280, (2007).
35. M. Miyayama, T. Nagamoto and O. Omoto, *Thin Solid Films*, **300**, 299 (1997).
36. Y. Noguchi, M. Miyayama, K. Oikawa, T. Kamiyama, M. Osada and M. kakahana, *Jpn. J. Appl. Phys.* **41**, 7062 (2002).
37. W. Wang, J. Zhu, X. Y. Mao and X. B. Chen, *Mater. Res. Bull.*, **42**, 274 (2007)
38. S. Takahashi & M. Takahashi. *Jpn. J. Appl. Phys.*, **11**, 31 (1972).
39. Y. Noguchi, I Miwa, Y.Goshima & M. Miyayama *Jpn. J. Appl. Phys.*, **39**, L1259 (2000).